

Space performance of the multistage labyrinthine SMEI baffle

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ABSTRACT

The Solar Mass Ejection Imager (SMEI) was launched on 6 January 2003, and shortly thereafter raised to a nearly circular orbit at 840 km. Three SMEI CCD cameras on the zenith-nadir oriented CORIOLIS spacecraft cover most of the sky beyond about 20° from the Sun, each 102-minute orbit. Data from this instrument provide precision visible-light photometric sky maps. Once starlight and other constant or slowly varying backgrounds are subtracted, the residue is mostly sunlight that has been Thomson-scattered from heliospheric electrons. These maps enable 3-dimensional tomographic reconstruction of heliospheric density and velocity. This analysis requires 0.1% photometry and background-light reduction below one S10 (the brightness equivalent of a 10th magnitude star per square degree). Thus 10^{-15} of surface-brightness reduction is required relative to the solar disk. The SMEI labyrinthine baffle provides roughly 10^{-10} of this reduction; the subsequent optics system provides the remainder. We analyze data obtained over two years in space, and evaluate the full system's stray-light rejection performance.

Keywords: Background-light reduction; space optics; very wide-angle optical designs

1. INTRODUCTION

The CORIOLIS spacecraft, with the Solar Mass Ejection Imager (SMEI)¹⁻⁶ onboard, was launched on 6 January 2003. Shortly thereafter it was raised to a nearly circular, polar terminator orbit at 840 km. Three SMEI cameras on the zenith-nadir oriented spacecraft view most of the sky beyond about 20° from the Sun. Each camera consists of a labyrinthine baffle protecting the optical pupil from stray light, plus a two-mirror optical system to place the 60° × 3° field of view upon a CCD detector. The three cameras are aligned to view together a 160° × 3° strip oriented roughly perpendicular to the satellite's orbital plane, thus providing a nearly complete visible-light sky map each 102-minute orbit. SMEI was designed and constructed by a team of scientists and engineers from the U.S. Air Force Research Laboratory, UCSD, the University of Birmingham UK, Boston College, and Boston University. Financial support is being provided by the Air Force, the University of Birmingham and NASA.

A previous article presents calculations and laboratory measurements of the SMEI baffles⁷. To enable photometric measurements of heliospheric structures such as coronal mass ejections, each camera's combination of baffle and optics must meet the stray-light surface-brightness reduction specification of 10^{-15} relative to the solar disk. By design, the baffle provides roughly 10^{-10} , and the subsequent optics an additional 10^{-5} . Now that in-space SMEI data are available, end-to-end performance can be evaluated. This article presents results of a stray-light analysis of flight data spanning a two years period of performance.

2. BAFFLE AND OPTICS

The SMEI baffle is described in detail elsewhere⁷. Figures 1 and 2 show the baffle layout and summarize its stray-light reduction, for light exiting the baffle's final aperture Z0, as a function of the incident angles[†]. However, full end-to-end system performance also includes further stray-light rejection provided by the subsequent optics.

[†] The incident angles $[\Theta_x, \Theta_y]$ used here and in the previous article⁷ represent the intersection of two planes rotated respectively about the Y- and X-axes; the camera points towards the Z-axis and the long direction of its field of view lies along the X-axis. Θ_x (or Θ_y) is the angle between the line of intersection and the Y-Z (or X-Z) plane. The corresponding direction cosines $[C_x, C_y]$ are, respectively, $C_x = \sin(\Theta_x) \times (1 + \tan^2(\Theta_y) \cos^2(\Theta_x))^{-1/2}$ and $C_y = \sin(\Theta_y) \times (1 + \tan^2(\Theta_x) \cos^2(\Theta_y))^{-1/2}$. Then, as usual, $C_z = (1 - C_x^2 - C_y^2)^{1/2}$. This latter is the obliquity factor mentioned in §7 of the previous article: the formula given there is correct only in lowest order.

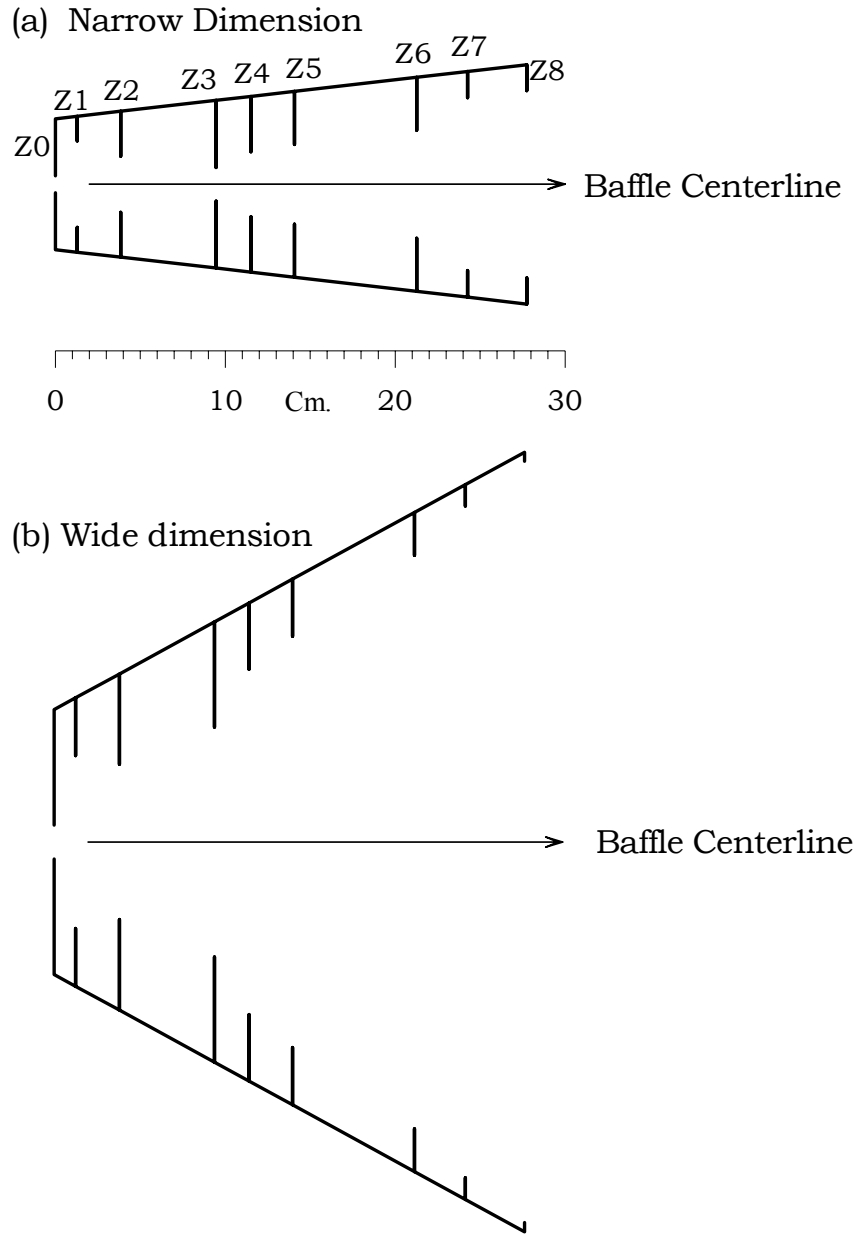


Figure 1. SMEI baffle design. Aperture numbers Z0 to Z8 advance with distance towards incident light along the baffle centerline. The strategic primary apertures are Z0 (rear of 1st stage and pupil of the subsequent SMEI optics, located out of sight off to the left here), Z3 (2nd stage rear), Z6 (3rd stage rear), and finally Z8 (3rd stage front). In the wide dimension the Z3, Z6 and Z8 edges line up, so here the “2nd stage” of the baffle extends from Z3 to Z8 and includes Z6. Secondary apertures Z4 → Z7 are placed between Z3 and Z8 to block illumination and/or viewing of the septum bottoms, as described in an earlier article⁷. All surfaces here are coated with Martin Black⁸.

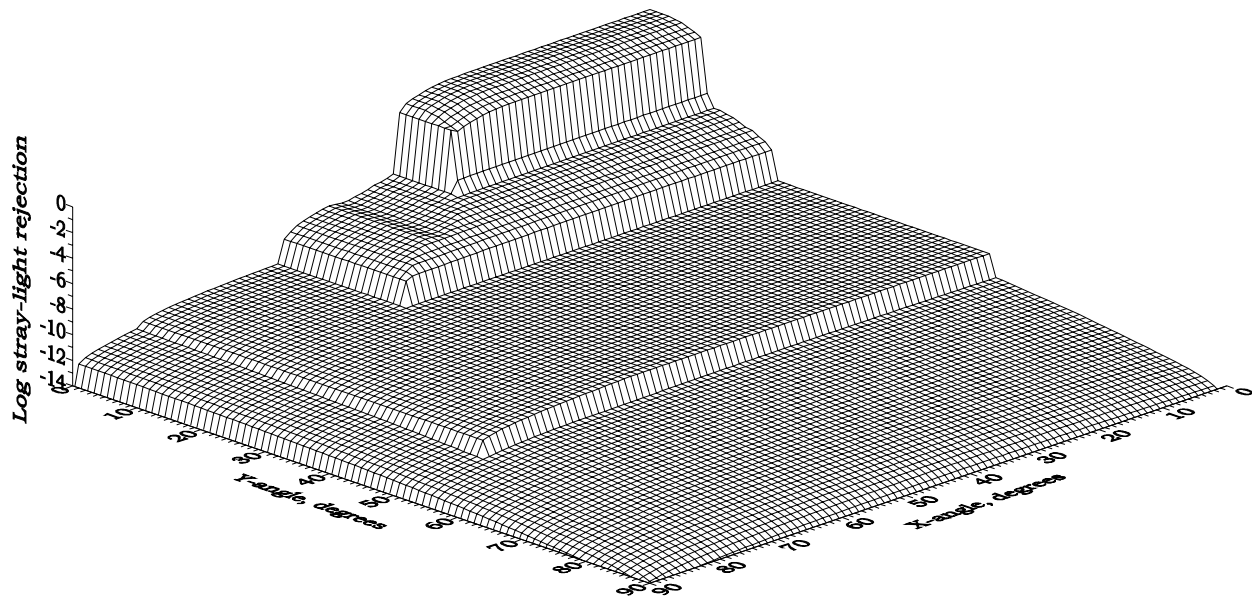


Figure 2. Plot of SMEI baffle stray-light rejection versus incident-light angles, normalized to unity at [0,0], from reference 7.

Figure 3 shows an exploded view of the optical system and figure 4 is a cross-section showing light rays incident at $\Theta_y = 0^\circ$. In figure 4 light at the center of the field of view (FOV) enters through the final baffle aperture Z0. It then reflects from mirrors M1 and M2 and downward to the CCD. Vane V serves some of the function of coronagraphic Lyot and field stops by blocking light at wide angles or from outside of the FOV. The light-sensitive portion of the CCD is the left-hand rectangle in figure 3, whereas the covered (frame-transfer) portion is the right-hand rectangle. These optical components, except the aperture and CCD, are all mounted to the base plate shown, and this portion of the optics is mounted to the baffle and entirely enclosed by a strongbox structure. The SMEI optics design is an off-axis Newtonian telescope as viewed in the Y-Z plane (figure 4), combined with a Schmidt camera without its corrector plate, as viewed in the X-Z plane. Mirror M1 was made from a 6061-T6 aluminum ring, diamond-turned to the correct figure, which was then separated into quadrants. To avoid stress-relief figure change, the quadrants were separated prior to the final diamond turning. The M2 mirrors were made two at a time in half-circles. To avoid increasing wide-angle scattered light, all mirrors were not “post-polished”. The optical surface was pure-aluminum-flashed and overcoated with SiO_x . Because of the major concern here about stray light, the mirror specification is a surface-brightness reduction of at least 10^{-4} , for light scattered by more than 1° . The stray-light distribution of individual flight mirrors was measured using a photodiode and illuminating a portion of the mirror with a laser. The above specification was confirmed for light scattering in the direction perpendicular to the residual diamond-turning grooves. In the direction along these grooves, the intensity of scattered light was even smaller, by about a factor of ten. The mirrors were manufactured by Hyperfine Inc., Boulder CO⁹. More detail about the optics is provided in the article by Eyles *et al.*⁵. The SMEI mission is described in Jackson *et al.*⁶.

3. STRAY LIGHT -- REDUCTION IN THE OPTICS BOX

When light enters through Z0 at angles outside the SMEI photometric FOV (*i.e.* when $|\Theta_x| > 30^\circ$, or $\Theta_y < -1.8^\circ$, or $\Theta_y > +1.5^\circ$), only a portion of it reaches the CCD. When the incident light lies *just* outside the FOV, the view of the sky is vignetted. Figure 5 illustrates two cases for incident light just outside of the FOV. When $+2^\circ < \Theta_y < +4.2^\circ$ a portion of the light reaches the CCD having reflected only from mirror M1 and not M2 (figure 5b). In this case the light does not enter the photometric FOV, but it often floods the CCD if a bright object such as the Moon is nearby. Data with the Moon inside $-3^\circ < \Theta_y < +4.5^\circ$ saturate the CCD and must be discarded. Out to $|\Theta_y| < 6.5^\circ$ the Moon still illuminates a portion of M1 and usually causes a shift in the electronic offset which is difficult to model and remove. A similar but smaller offset shift occurs when Venus lies within the FOV. These data frames are usually also discarded.

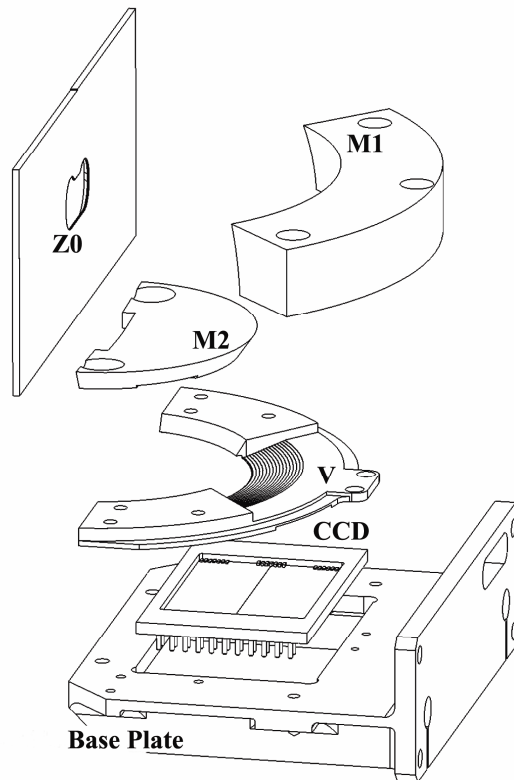


Figure 3. Exploded view of the SMEI optics. See also figures 4 and 5. See text for descriptions of the various items shown here.

The reflective part of mirror M1 is not illuminated by sky objects having $|\Theta_y| > 6.5^\circ$. In this case, light entering through Z0 illuminates the black-anodized interior of the optics box, and can reach the CCD when re-radiated from certain places on the optics-box top or sides. The CCD FOV subtends typically about 0.01 steradians from these locations. Hence light reflected and reaching the CCD FOV is reduced by a factor of 20 for the reflectivity, another factor of 300 for the solid angle, and finally is spread out over the equivalent of roughly 200 square degrees of sky, for an expected total surface-brightness reduction of about 10^{-6} . Thus, for sunlight inside of the limits $|\Theta_x| < 58^\circ$ and $|\Theta_y| < 19.5^\circ$ (the second level from the top in figure 1 above, with baffle rejection between 10^{-6} and 10^{-7}), the baffle+optics rejection is only about 10^{-12} . This would yield a stray-light sky background of between 10^2 and 10^3 S10 units (defined above in the abstract). This background is likely to be unacceptable for 3-dimensional tomographic reconstructions but might be adequate for some 2-dimensional image analysis. However, we show below that excessive light outside of the FOV ruins these CCD frames, thus rendering moot such data recovery.

Light just outside of the FOV can also interfere with precision photometry if a path within the optics box scatters it into the photometric FOV. The series of steps, placed as shown below in figure 3 on vane V's gradual sloped surface nearest mirror M2 and above the CCD, remove background light. Had this surface been smooth, this light would scatter into the photometric FOV by grazing-angle reflection followed by a reflection from M2. Similarly, the "walls" terminating these steps on either end of the stepped section scatter light which in their absence (and a continuation around of the series of steps) would be focused safely outside of the photometric FOV. Some of this scattered light reaches the photometric FOV directly following the scattering and another portion after reflecting from M2. This is a small but significant stray-light source (here called the "glare") for the two cameras viewing closest to the Sun. A subtraction procedure, still under development, yields promising preliminary results and will likely restore full SMEI performance. Strictly speaking, *this* stray-light source is not due to the baffle itself and we thus defer detailed discussion to a future article.

4. RESULTS WITH SMEI FLIGHT DATA

The SMEI camera doors were opened at about 2000h GMT on 1 February 2003. About two dozen orbits of “contingency” data were taken with the various onboard-binning modes. Then roughly one orbit for each camera was recorded and the full 256×1272 pixel CCD readout sent down through telemetry. This particular calibration sequence has been repeated annually, and the results presented here are entirely from these data.

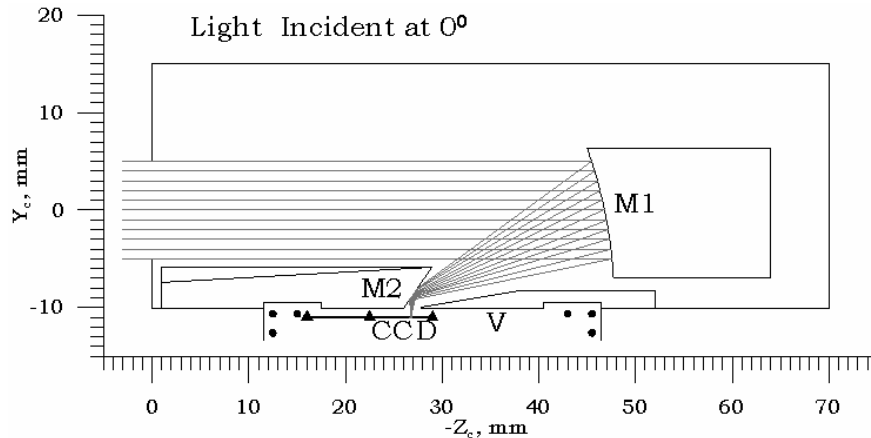


Figure 4. Elevation view of the SMEI optics, a cross-section through the center of figure 3, but with the optical elements now in proper relationship to one another. The baffle structure is off to the left. Corners of the CCD frame are shown as round dots, and the light-sensitive portion of the CCD (the left-hand rectangle in figure 3) is marked by three triangles. Mirrors M1 and M2 are diamond-turned aluminum. The actual top of the optics box is at $Y_c \approx 85$ mm, much higher than shown here; this provides room for the shutter and a larger volume for disposing of unwanted stray light. See figure 9 in reference 5 for a more detailed picture showing the shutter.

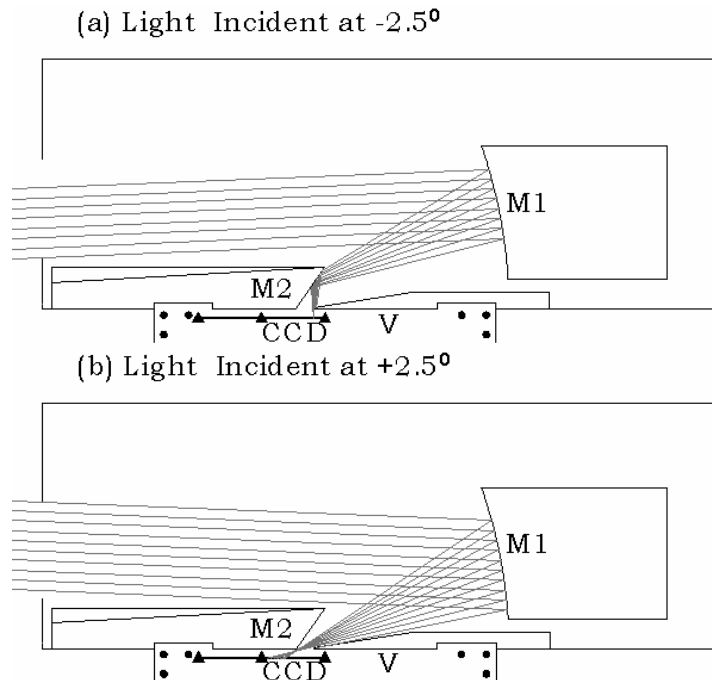


Figure 5. View as in figure 4, but with axes suppressed and light incident at $\Theta_y = \pm 2.5^\circ$. In (a), rays reflecting from M1 but hitting the top of M2 are not shown. In (b), the lowest ray, also not shown, intercepts the top of M2 before reaching M1. In each case vane V intercepts some of the light. In case (b) the light reaching the CCD is outside of the photometric FOV.

Figures 6 and 7 respectively show three data frames each from camera 3, which views closest to the Sun, and from camera 2. These data include the entire available CCD rectangle rather than the usual “region of interest”, a ~100 pixel wide circular band extending the full width of the CCD. These figures illustrate most SMEI stray-light features, and since they span two years in space, also place stringent limits on possible degradation of the Martin-black coating which, together with the baffle geometry, control the amount of stray light. Only a portion of the above circular band is suitable as the *photometric* FOV, both due to vignetting of the instrument aperture Z0 by other optical elements as seen in figure 5, and from the stray-light bands seen in both of these pictures. The photometric FOV is a ~60-pixel wide band roughly centered within the region of interest. This latter band contains the unvignetted view of the sky and also is mostly free from stray light.

The brightness unit here is an analog-to-digital unit (ADU), one of which on the CCD is 4.7 electrons detected by a single pixel, each of which in turn covers on average a $0.05^\circ \times 0.05^\circ$ patch on the sky⁵. In figure 6, various stars are superimposed on a smoother background varying from about 150 ADUs at the end away from the Sun, to about 350 ADUs at the end towards it. The corresponding sky surface brightnesses in these areas are about 400 and 900 S10 units respectively. Roughly twice-brighter bands of light occur both above and below this photometric FOV. The lower one is sharply cut off by the close edge of vane V, while the upper is more diffuse. These are out-of-focus images of the long edges of the baffle’s aperture Z3. This aperture is illuminated by sunlight entering the front of the baffle, which scatters several times before finding its way to the rear of Z3, and then scatters once more through aperture Z0 to reach the SMEI optics and the CCD. The SMEI optics design places the out-of-focus images of these aperture Z3 vane edges safely just outside of the photoelectric FOV. The “glare” mentioned above at the end of section 3 is a hard-to-see 30 ADUs broad blob centered at about [100, 200] pixels here, and another half as bright at [1200, 200].

In figure 6, the Sun is at $[\Theta_x, \Theta_y] \approx [36^\circ, 35^\circ]$, so the baffle-only performance (figure 2) for these pictures lies upon the 3rd level down from the top at $[0^\circ, 0^\circ]$ and has an expected light reduction of 10^{-10} . If *all* of this stray light were to reach the CCD, it would total about 10^5 S10 units: here 10^8 electrons or 2.5×10^7 ADUs. The inner and outer bands in figure 6a respectively contain about 7.4×10^6 and 6.5×10^6 ADUs. Thus in this case about half of the stray sunlight entering the optics box does reach the CCD, but is largely confined to these two bands, as expected from the original design calculations.

In the pictures of figure 7, camera 2’s axis is at $[\Theta_x, \Theta_y] \approx [87^\circ, 30^\circ]$, see figure 2. Baffle-only rejection at this location is about 10^{-12} , so the stray light in the two bands should be reduced about 100-fold; indeed, the contents of the bands are here reduced to about 2×10^5 ADUs each. The stray light is mostly confined to the end of the FOV opposite to the Sun. This results from the sunlight initially illuminating only a small part of the outermost septum bottom in the baffle’s wide dimension (see figure 1b). From there it scatters once more, toward the opposite side, to pass through the Z3 aperture and then a second time (reversing its direction) to illuminate the Z3 rear surface. Here geometry favors illuminating that end of Z3 nearest the original illumination.

Figure 8 shows the combined data for an entire camera 3 orbit, excluding two portions where the CCD response saturated due to well overfilling, a large one due to the Sun, and a smaller one due to the Moon. An onset of the Sun’s exclusion occurs abruptly when this has $|\Theta_y| < 19.5^\circ$, exactly where expected from the design calculations and earlier calibrations⁷. In this case a factor of 10^4 more light floods through the critical aperture Z3 into the baffle’s innermost stage, and much of it passes through Z0 into the optics box. In this case, most of the CCD wells become overfilled, rendering the data frame unusable. To create this figure, the CCD pixels within the photometric FOV for each data frame were transformed into Right Ascension (R.A.) and Declination (δ) coordinates in a hierarchical-triangle spherical-coordinate system⁶, using the spacecraft orientation quaternions and known camera orientations⁵. The observed SMEI surface brightness (ADUs per pixel) is first corrected for geometrical effects, and then mapped to this coordinate system. Here, for each hierarchical triangle lying within a pixel, the surface brightness is averaged with contributions from the other frames and the result then appropriately binned for the figure.

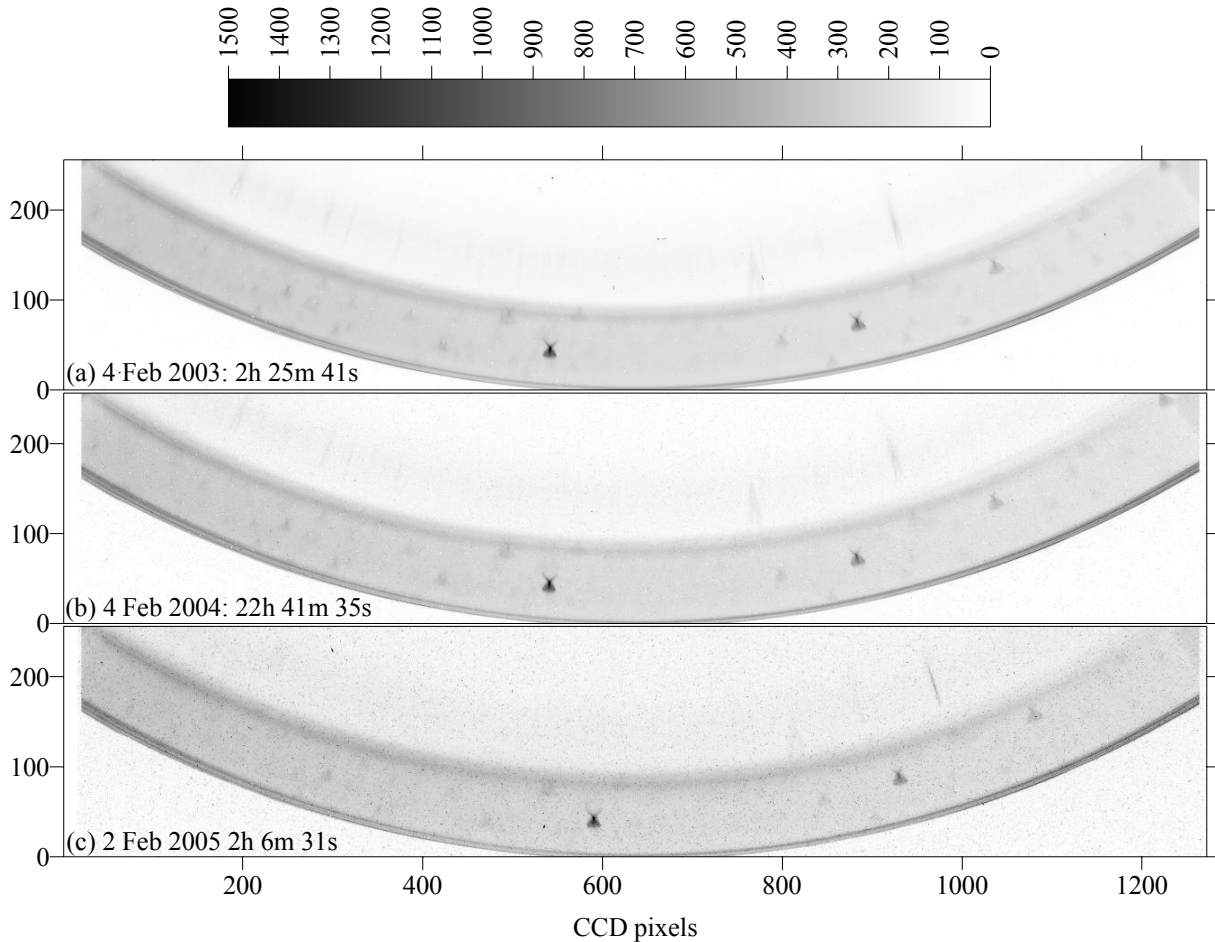


Figure 6. Three data frames at roughly one-year intervals from SMEI camera 3, the one viewing closest to the Sun. The intensity scale is ADUs, where 1 ADU is 4.7 electrons in a pixel. In the 4 second exposure, a 10th magnitude star has about 800 electrons detected in the CCD. In this figure, an electronic offset pedestal, dark current and a scaled “hot pixel” pattern are removed; individual-pixel impacts of Earth radiation-belt particles and cosmic rays have been removed by a simple spatial filter. The camera views a roughly 60° long by 3° high rectangle in the sky. In this particular figure the dimension along the arc is roughly parallel to lines of constant R.A. These pictures are left-right reversed from actual sky, and the sky’s rectangular FOV is mapped onto this curved shape by the above-described SMEI optics. The camera axis for the middle frame points towards right ascension R.A. = 23h 24m and declination $\delta = +15.5^\circ$. The bright star near the center is α Peg, another to the right is γ Peg. At the time of these pictures, the Sun was at R.A. = 21h 11m and $\delta = -16.2^\circ$, about 600 pixels below the origin of coordinates above. Several faint distorted star images are visible also above and within the bright arc, coming from stars just north of this FOV, whose light reaches the CCD by slipping under mirror M2 by the path shown in figure 5b. Noise “aging” in camera 3 is visible as a background of mostly isolated spots whose density is increasing from (a) to (c).

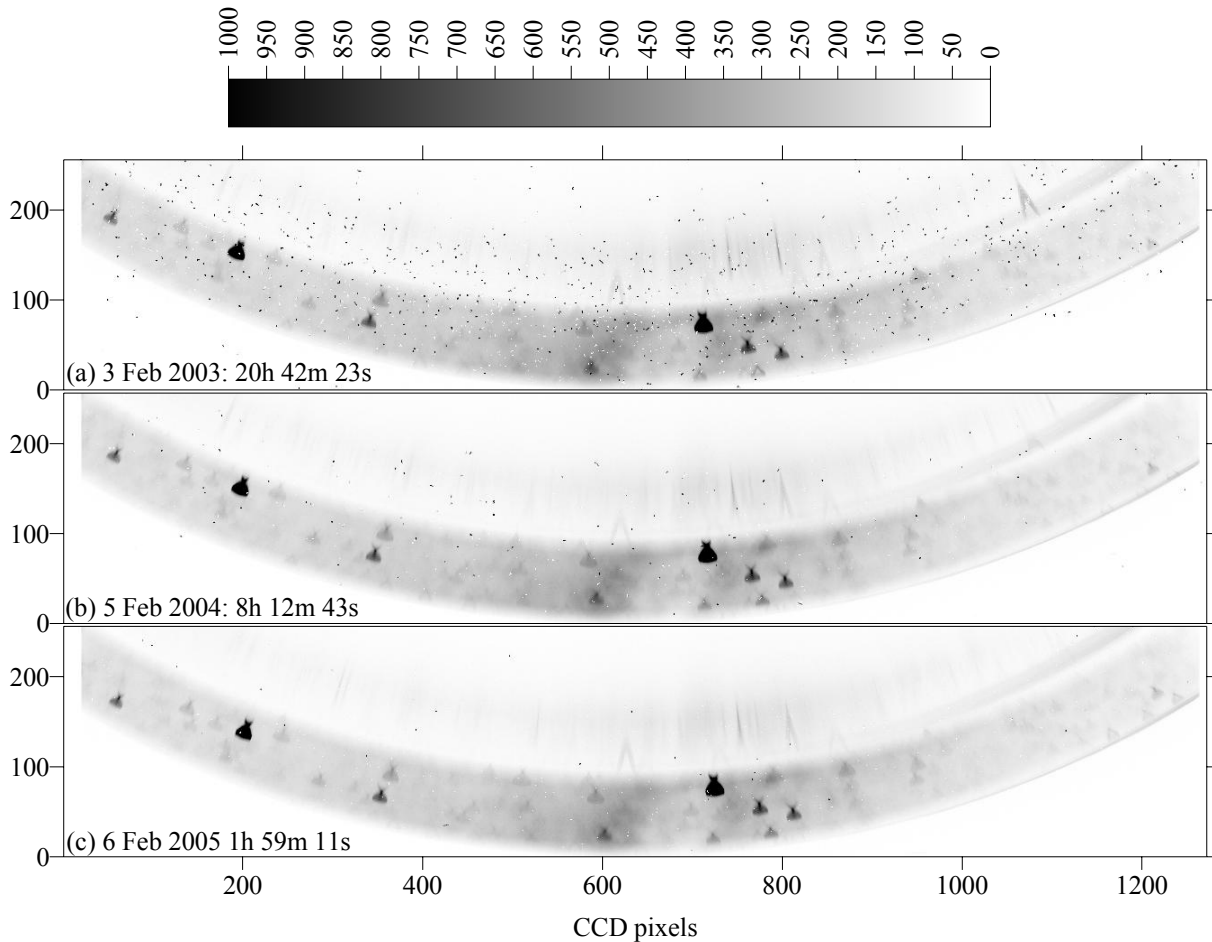


Figure 7. Three frames at roughly one-year intervals from SMEI camera 2, the one viewing about 90° from the Sun. The intensity scale is ADUs, note the brightness-scale adjustment relative to figure 6. Here the stray-light bands are visible only on the right hand end of the picture, and even there with intensity significantly reduced from figure 6. The camera axis for these pictures points towards R.A. = 13h 6m and declination $\delta = -64.7^\circ$. This southerly viewing sometimes places the SMEI spacecraft near the South Atlantic Anomaly, which can contribute significant particle-hit background on the CCD, here especially in (a). Aging in camera 2 is much less visible here than in camera 3 (figure 6) because the CCD in camera 2 operates at a significantly colder temperature.

5. PERFORMANCE STABILITY OVER TWO YEARS

The annual calibration procedure enables tracking the stray-light performance of SMEI, at present through two full years of operation. If the Martin-black coatings degrade with time, the amount of light in the two bands outside of the photometric FOV will increase in individual data frames like those shown in figures 6 and 7. In figure 6, the contents of the bands increases by about 0.5×10^6 ADUs per year. If this increase is due to an increase in the reflectivity of Martin black due to contamination or some other temporal change, and if the reflectivity increase is the same factor for each of the ~ 4 surfaces encountered before this light can pass through aperture Z0, then a single surface's reflectivity increases for each by a factor of only 1.02 per year. The lower intensity of the bands in figure 7, plus variation in the contribution from stars close to or occupying the bands precludes a definitive separate measurement of baffle reflectivity change with time for camera 2.

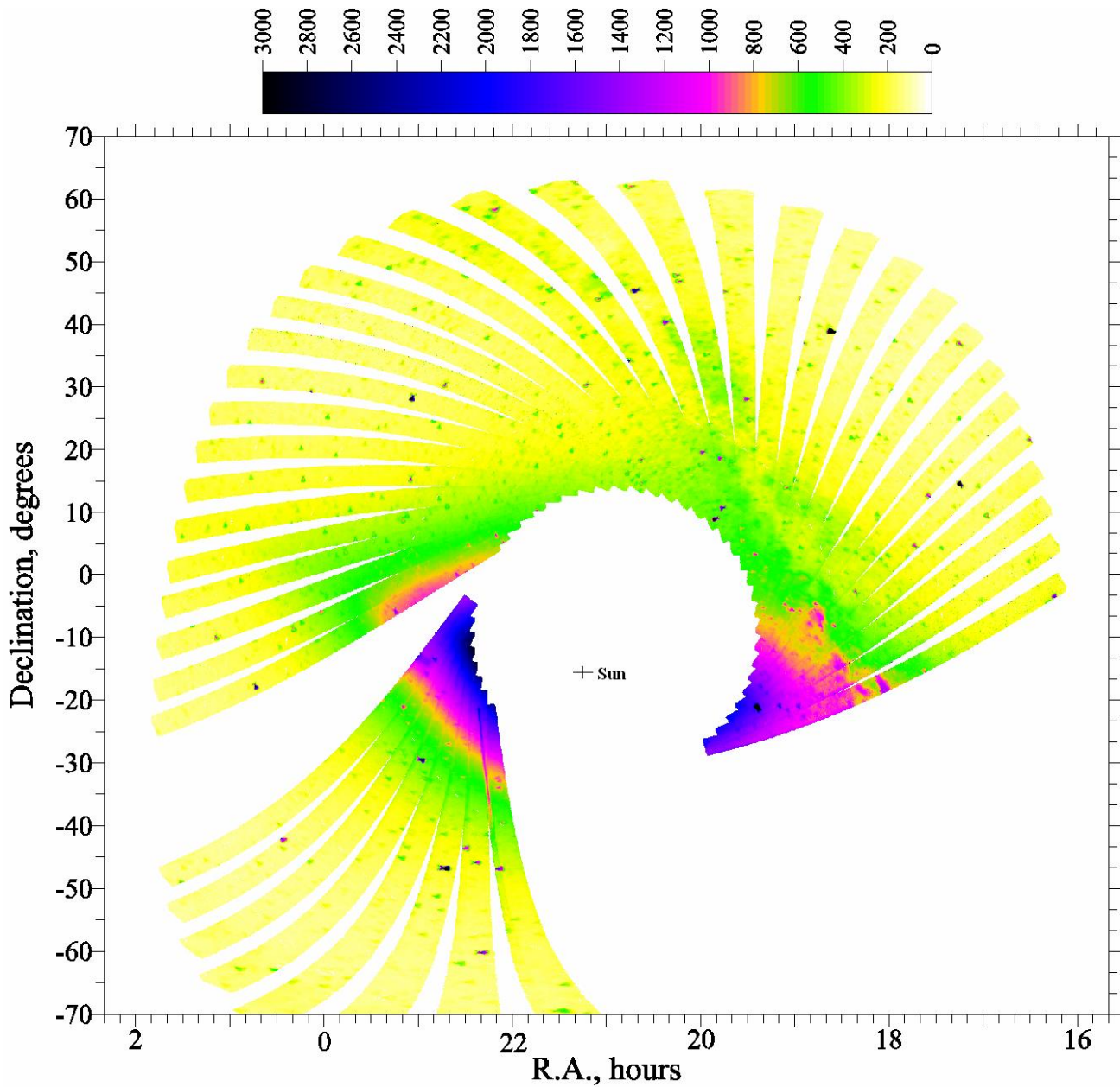


Figure 8. A composite of a sequence of frames like figure 6, recorded on 4 February 2003, covering one full orbit between 2h 14m 13s and 3h 45m 49s. A small cross marks the Sun's location at this time. Data are excluded between 3h 7m 29s and 3h 36m 37s (the large gap to the lower right in the otherwise nearly full circular pattern) when the Sun has $\Theta_y < 19.5^\circ$ and thus bright stray sunlight overfills the CCD wells. The frames come in groups of typically six separated by 4 seconds each, and then with a time gap of about 1.5 minutes between groups while data are transferred to the spacecraft onboard buffer. Another smaller gap, to the lower left, is caused by well overfilling from the Moon at R.A. = 23h 19.5m and $\delta = -9.9^\circ$. Some scattered light is visible nearby this, and nearby Venus which here is at R.A. = 17h 58m, $\delta = -21.4^\circ$. Also, Mercury is visible as a bright object at R.A. = 19h 25m, $\delta = -21.2^\circ$.

6. DISCUSSION AND CONCLUSIONS

The final definitive verification of the full 10^{-15} end-to-end performance of the SMEI baffle and optics combination awaits a detailed modeling of the zodiacal light and the demonstration of a successful subtraction both of this and of background starlight from SMEI sky maps similar to figure 8. However, figures 6 and 7 both do show that the stray background light is observed roughly where expected by the original design, and in the right amount. Light intensity drops off rapidly as one moves from either of the two bands along the narrow dimension of the band, radially towards and into the photometric FOV. This begins about 1° away from the bands, although some of the light might be picked up were the photometric FOV not carefully placed. Roughly a factor of 100 dropoff relative to the band's surface brightness is required to meet the SMEI specification. If the main mechanism for light in these bands reaching the photometric FOV were just scattering from mirror M1, the individual-mirror scattering measurements indicate that even better performance than this would be realized. The "glare" discussed above was unanticipated in the original design, but in a future SMEI-type instrument can be reduced by texturing the walls terminating the stepped portion of vane V, and/or covering reflective portions at both ends of mirror M1 that are not illuminated by light legitimately entering the photometric FOV. If the subtraction procedure for this described above in Section 3 is fully successful, and the verification described at the beginning of this section is thus realized, then SMEI indeed will have achieved its original stray-light surface-brightness reduction specification of 10^{-15} relative to the Sun.

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